5.1: Sturm-Liouville problems

5.1.1 Boundary value problems

We have encountered several different eigenvalue problems such as:

\[ X''(x) + \lambda X(x) = 0 \]

with different boundary conditions

\[ X(0) = 0 \quad X(L) = 0 \quad \text{(Dirichlet or)} \]
\[ X'(0) = 0 \quad X'(L) = 0 \quad \text{(Neumann or)} \]
\[ X'(0) = 0 \quad X(L) = 0 \quad \text{(Mixed or)} \]
\[ X(0) = 0 \quad X'(L) = 0 \quad \text{(Mixed)} \]

For example for the insulated wire, Dirichlet conditions correspond to applying a zero temperature at the ends, Neumann means insulating the ends, etc.… Other types of endpoint conditions also arise naturally, such as the Robin boundary conditions

\[ hX(0) - X'(0) = 0 \quad hX(L) + X'(L) = 0, \]

for some constant \( h \). These conditions come up when the ends are immersed in some medium.
Boundary problems came up in the study of the heat equation \(u_t=ku_{xx}\) when we were trying to solve the equation by the method of separation of variables. In the computation we encountered a certain eigenvalue problem and found the eigenfunctions \(X_n(x)\). We then found the eigenfunction decomposition of the initial temperature \(f(x)=u(x,0)\) in terms of the eigenfunctions

\[
[f(x)= \sum_{n=1}^{\text{infty}} c_n X_n(x)].
\]

Once we had this decomposition and found suitable \(T_n(t)\) such that \(T_n(0)=1\) and \(T_n(t)X(x)\) were solutions, the solution to the original problem including the initial condition could be written as

\[
[u(x,t)= \sum_{n=1}^{\text{infty}} c_n T_n(t)X_n(x)].
\]

We will try to solve more general problems using this method. First, we will study second order linear equations of the form

\[
\frac{d}{dx}\left( p(x)\frac{dy}{dx} \right)-q(x)y+\lambda r(x)y=0.
\]

Essentially any second order linear equation of the form \(a(x)y''+b(x)y'+c(x)y+\lambda d(x)y=0\) can be written as (5.1.5) after multiplying by a proper factor.

Example \(\PageIndex{1}\): Sturm-Liouville Problem

Put the following equation into the form (5.1.5):

\[
x^2y''+xy'+(\lambda x^2-n^2)y=0.
\]

Multiply both sides by \(\frac{1}{x}\) to obtain

\[
\frac{1}{x}(x^2y''+xy'+(\lambda x^2-n^2)y)=xy''+y'+ \left( \lambda x -\frac{n^2}{x}\right)y= \frac{d}{dx}\left( x \frac{dy}{dx} \right)-\frac{n^2}{x}y+\lambda xy=0.
\]

The so-called Sturm-Liouville problem is to seek nontrivial solutions to

\[
\frac{d}{dx}\left( p(x)\frac{dy}{dx} \right)-q(x)y+\lambda r(x)y=0, \quad a<x<b, \quad \alpha_1y(a)-\alpha_2y'(a)=0, \quad \beta_1y(b)+\beta_2y'(b)=0.
\]

In particular, we seek \(\lambda\)s that allow for nontrivial solutions. The \(\lambda\)s that admit nontrivial solutions are called the eigenvalues and the corresponding nontrivial solutions are called eigenfunctions. The constants \(\alpha_1\) and \(\beta_1\) should not be both zero, same for \(\beta_2\).

Theorem 5.1.1. Suppose \(p(x), p'(x),q(x)\) and \(r(x)\) are continuous on \([a,b]\) and suppose \(p(x)>0\) and \(r(x)>0\) for all \(x\) in \((a,b)\). Then the Sturm-Liouville problem (5.1.8) has an increasing sequence of eigenvalues

\[
[\lambda_1<\lambda_2<\lambda_3<\cdots]
\]

such that
\[ \lim_{n \to \infty} \lambda_n = +\infty \]

and such that to each \( \lambda_n \) there is (up to a constant multiple) a single eigenfunction \( y_n(x) \).

Moreover, if \( q(x) \geq 0 \) and \( \alpha_1, \alpha_2, \beta_1, \beta_2 \geq 0 \), then \( \lambda_n \geq 0 \) for all \( n \).

Problems satisfying the hypothesis of the theorem are called regular Sturm-Liouville problems and we will only consider such problems here. That is, a regular problem is one where \( p(x), q(x) \) and \( r(x) \) are continuous, \( p(x) > 0 \), \( r(x) > 0 \), \( q(x) \geq 0 \), and \( \alpha_1, \alpha_2, \beta_1, \beta_2 \geq 0 \). Note: Be careful about the signs. Also be careful about the inequalities for \( r(x) \) and \( p(x) \), they must be strict for all \( x \)!

When zero is an eigenvalue, we usually start labeling the eigenvalues at \( 0 \) rather than at \( 1 \) for convenience.

Example \( \PageIndex{2} \):

The problem \( y'' + \lambda y, 0 < x < L, y(0) = 0, y(L) = 0 \) is a regular Sturm-Liouville problem. \( p(x) = 1, q(x) = 0, r(x) = 1 \), and we have \( p(x) > 0 \) and \( r(x) > 0 \). The eigenvalues are \( \lambda_n = \frac{n^2 \pi^2}{L^2} \) and eigenfunctions are \( y_n(x) = \sin\left(\frac{n \pi}{L} x\right) \). All eigenvalues are nonnegative as predicted by the theorem.

Exercise \( \PageIndex{1} \):

Find eigenvalues and eigenfunctions for

\[ y'' + \lambda y = 0, \quad y'(0) = 0, \quad y'(1) = 0. \]

Identify the \( p, q, r, \alpha_j, \beta_j \). Can you use the theorem to make the search for eigenvalues easier? (Hint: Consider the condition \( -y'(0) = 0 \))

Example \( \PageIndex{3} \):

Find eigenvalues and eigenfunctions of the problem

\[ y'' + \lambda y = 0, \quad 0 < x < 1, \quad hy(0) - y'(0) = 0, \quad y'(1) = 0, \quad h > 0. \]

These equations give a regular Sturm-Liouville problem.

Exercise \( \PageIndex{2} \):

Identify \( p, q, r, \alpha_j, \beta_j \) in the example above.

First note that \( \lambda \geq 0 \) by Theorem 5.1.1. Therefore, the general solution (without boundary conditions) is

\[ y(x) = A \cos(\sqrt{\lambda} x) + B \sin(\sqrt{\lambda} x) \quad \text{if} \quad \lambda > 0, \]

\[ y(x) = Ax + B \quad \text{if} \quad \lambda = 0. \]
Let us see if \( \lambda=0 \) is an eigenvalue: We must satisfy \( 0=hB-A \) and \( A=0 \), hence \( B=0 \) (as \( h>0 \)), therefore, \( \lambda=0 \) is not an eigenvalue (no nonzero solution, so no eigenfunction).

Now let us try \( h>0 \). We plug in the boundary conditions.

\[
0 = hA - \sqrt{\lambda}B, \quad 0 = -A\sqrt{\lambda}\sin(\sqrt{\lambda}) + B\sqrt{\lambda}\cos(\sqrt{\lambda}).
\]

If \( A=0 \), then \( B=0 \) and vice-versa, hence both are nonzero. So \( B=\frac{hA}{\sqrt{\lambda}} \), and \( 0=-A\sqrt{\lambda}\sin(\sqrt{\lambda})+\frac{hA}{\sqrt{\lambda}}\sqrt{\lambda}\cos(\sqrt{\lambda}) \). As \( A \neq 0 \) we get

\[
0 = -\sqrt{\lambda}\sin(\sqrt{\lambda}) + h\cos(\sqrt{\lambda}),
\]

or

\[
\left[ \sqrt{\lambda} \right] = \tan \left[ \sqrt{\lambda} \right].
\]

Now use a computer to find \( \lambda_n \). There are tables available, though using a computer or a graphing calculator is far more convenient nowadays. Easiest method is to plot the functions \( \frac{h}{x} \) and \( \tan(x) \) and see for which

\( x \) they intersect. There is an infinite number of intersections. Denote by \( \sqrt{\lambda_1} \) the first intersection, by \( \sqrt{\lambda_2} \) the second intersection, etc. For example, when \( h=1 \), we get that \( \sqrt{\lambda_1} \approx 0.86, \sqrt{\lambda_2} \approx 3.43, ... \). That is \( y_1 \approx 0.74, y_2 \approx 11.73, ... \). A plot for \( h=1 \) is given in Figure 5.1. The appropriate eigenfunction (let \( A=1 \) for convenience, then \( B=\frac{h}{\sqrt{\lambda}} \)) is

\[
y_n(x) = \cos(\sqrt{\lambda_n}x) + \frac{h}{\sqrt{\lambda_n}}\sin(\sqrt{\lambda_n}x).
\]

When \( h=1 \) we get (approximately)

\[
y_1(x) \approx \cos(0.86x) + \frac{1}{0.86} \sin(0.86x), \quad y_2(x) \approx \cos(3.43x) + \frac{1}{3.43} \sin(3.43x), ....
\]
5.1.2 Orthogonality

We have seen the notion of orthogonality before. For example, we have shown that \( \sin(nx) \) are orthogonal for distinct \( n \) on \( [0, \pi] \). For general Sturm-Liouville problems we will need a more general setup. Let \( r(x) \) be a weight function (any function, though generally we will assume it is positive) on \( [a, b] \). Two functions \( f(x) \), \( g(x) \) are said to be orthogonal with respect to the weight function \( r(x) \) when

\[
\int_a^b f(x)g(x)r(x)dx = 0.
\]

In this setting, we define the inner product as

\[
\langle f, g \rangle \overset{\text{def}}{=} \int_a^b f(x)g(x)r(x)dx,
\]

and then say \( \langle f \rangle \) and \( \langle g \rangle \) are orthogonal whenever \( \langle f, g \rangle = 0 \). The results and concepts are again analogous to finite dimensional linear algebra.

The idea of the given inner product is that those \( x \) where \( r(x) \) is greater have more weight. Nontrivial (nonconstant) \( r(x) \) arise naturally, for example from a change of variables. Hence, you could think of a change of variables such that \( d\xi = r(x)dx \).

We have the following orthogonality property of eigenfunctions of a regular Sturm-Liouville problem.

**Theorem 5.1.2.** Suppose we have a regular Sturm-Liouville problem

\[
\text{Figure 5.1: Plot of } \left( \frac{1}{x} \right) \text{ and } (\tan x).
\]
\[
\frac{d}{dx} \left( p(x) \frac{dy}{dx} \right) - q(x)y + \lambda r(x)y = 0, \quad \alpha_1 y(a) - \alpha_2 y'(a) = 0, \quad \beta_1 y(b) + \beta_2 y'(b) = 0.
\]

Let \(y_j\) and \(y_k\) be two distinct eigenfunctions for two distinct eigenvalues \(\lambda_j\) and \(\lambda_k\). Then
\[
\int_a^b y_j(x)y_k(x)r(x)dx = 0,
\]
that is, \(y_j\) and \(y_k\) are orthogonal with respect to the weight function \(r\).

Proof is very similar to the analogous theorem from § 4.1. It can also be found in many books including, for example, Edwards and Penney [EP].

\[\text{5.1.3 Fredholm alternative}\]

We also have the Fredholm alternative theorem we talked about before for all regular Sturm-Liouville problems. We state it here for completeness.

Theorem 5.1.3 (Fredholm alternative). Suppose that we have a regular Sturm-Liouville problem. Then either
\[
\frac{d}{dx} \left( p(x) \frac{dy}{dx} \right) - q(x)y + \lambda r(x)y = 0, \quad \alpha_1 y(a) - \alpha_2 y'(a) = 0, \quad \beta_1 y(b) + \beta_2 y'(b) = 0,
\]
has a nonzero solution, or
\[
\frac{d}{dx} \left( p(x) \frac{dy}{dx} \right) - q(x)y + \lambda r(x)y = f(x), \quad \alpha_1 y(a) - \alpha_2 y'(a) = 0, \quad \beta_1 y(b) + \beta_2 y'(b) = 0,
\]
has a unique solution for any \(f(x)\) continuous on \([a,b]\).

This theorem is used in much the same way as we did before in § 4.4. It is used when solving more general nonhomogeneous boundary value problems. The theorem does not help us solve the problem, but it tells us when a unique solution exists, so that we know when to spend time looking for it. To solve the problem we decompose \(f(x)\) and \(y(x)\) in terms of the eigenfunctions of the homogeneous problem, and then solve for the coefficients of the series for \(y(x)\).

\[\text{5.1.4 Eigenfunction series}\]

What we want to do with the eigenfunctions once we have them is to compute the eigenfunction decomposition of an arbitrary function \(f(x)\). That is, we wish to write
\[
f(x) = \sum_{n=1}^{\infty} c_n y_n(x),
\]
where \(y_n(x)\) the eigenfunctions. We wish to find out if we can represent any function \(f(x)\) in this way, and if so, we wish to calculate
\( c_n \)
(and of course we would want to know if the sum converges). OK, so imagine we could write \( \langle f(x) \rangle \) as (5.1.24). We will assume convergence and the ability to integrate the series term by term. Because of orthogonality we have

\[
\langle f, y_m \rangle = \int_a^b f(x)y_m(x)r(x)dx = \sum_{n=1}^{\infty} c_n \int_a^b f(x)y_n(x)r(x)dx = c_m \langle y_m, y_m \rangle .
\]

Hence,

\[
\langle y_m, y_m \rangle = \frac{\int_a^b f(x)y_m(x)r(x)dx}{\int_a^b (y_m(x))^2r(x)dx}.
\]

Note that \( \langle y_m \rangle \) are known up to a constant multiple, so we could have picked a scalar multiple of an eigenfunction such that \( \langle y_m, y_m \rangle = 1 \) (if we had an arbitrary eigenfunction \( \tilde{y}_m \), divide it by \( \sqrt{\langle \tilde{y}_m, \tilde{y}_m \rangle} \)). When \( \langle y_m, y_m \rangle = 1 \) we have the simpler form \( c_m = \langle f, y_m \rangle \) as we did for the Fourier series. The following theorem holds more generally, but the statement given is enough for our purposes.

Theorem 5.1.4. Suppose \( \langle f \rangle \) is a piecewise smooth continuous function on

\([a, b]\)

If \( \langle y_1, y_2, \ldots \rangle \) are the eigenfunctions of a regular Sturm-Liouville problem, then there exist real constants \( c_1, c_2, \ldots \) given by (5.1.26) such that (5.1.24) converges and holds for \( a < x < b \).

Example \( \PageIndex{4} \):

Take the simple Sturm-Liouville problem

\[
y'' + \lambda y = 0, \quad 0 < x < \frac{\pi}{2}, \quad y(0) = 0, \quad y' \left( \frac{\pi}{2} \right) = 0.\]

The above is a regular problem and furthermore we know by Theorem 5.1.1 that \( \langle \lambda \rangle \geq 0 \).

Suppose \( \langle \lambda = 0 \rangle \), then the general solution is \( y(x)Ax + B) \), we plug in the initial conditions to get \( y(0) = 0 \) and \( y' \left( \frac{\pi}{2} \right) = 0 \), hence \( \langle \lambda = 0 \rangle \) is not an eigenvalue. The general solution, therefore, is

\[ y(x) = A \cos(\sqrt{\lambda}x) + B \sin(\sqrt{\lambda}x) \].

Plugging in the boundary conditions we get \( y(0) = 0 \) and \( y' \left( \frac{\pi}{2} \right) = \sqrt{\lambda}B \). \( B \) cannot be zero and hence \( \langle \cos(\sqrt{\lambda}x) \rangle = 0 \). This means that \( \langle \sqrt{\lambda} \rangle \frac{\pi}{2} \rangle \) must be an odd integral multiple of \( \langle \frac{\pi}{2} \rangle \), i.e.

\[ \langle (2n-1) \frac{\pi}{2} \rangle = \sqrt{\lambda_n} \langle \frac{\pi}{2} \rangle \].

Hence

\[ \langle \lambda_n = (2n-1)^2 \rangle .\]

We can take \( \langle B = 1 \rangle \). Hence our eigenfunctions are
Finally we compute
\[ \int_0^{\frac{\pi}{2}} (\sin((2n-1)x))^2 dx = \frac{\pi}{4}. \]
So any piecewise smooth function on \([0, \frac{\pi}{2}]\) can be written as
\[ f(x) = \sum_{n=1}^{\infty} c_n \sin((2n-1)x), \]
where
\[ c_n = \frac{\langle f, y_n \rangle}{\langle y_n, y_n \rangle} = \frac{\int_0^{\frac{\pi}{2}} \sin((2n-1)x) dx}{\int_0^{\frac{\pi}{2}} (\sin((2n-1)x))^2 dx} = \frac{4}{\pi} \int_0^{\frac{\pi}{2}} f(x) \sin((2n-1)x) dx. \]
Note that the series converges to an odd \((2\pi)\)-periodic (not \((\pi)\)) \(\pi\)-periodic! extension of \(f(x)\).

Exercise \(\PageIndex{3}\)

In the above example, the function is defined on \((0 < x < \frac{\pi}{2})\), yet the series converges to an odd \((2\pi)\)-periodic extension of \((f(x))\). Find out how is the extension defined for \((\frac{\pi}{2} < x < \pi)\).

\(^1\)Named after the French mathematicians Jacques Charles François Sturm (1803–1855) and Joseph Liouville (1809–1882).

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